Performance Modeling of Ring Laser Gyro in Inertial Navigation System

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Abstract: In this paper, the performance of a Ring Laser Gyro based inertial navigation is investigated. Dynamic and stochastic modeling are applied to gyro simulation and performance evaluation. In the dynamic model, some parameters such as scale factor and environmental sensitivity have been determined, whereas in the stochastic model, the other parameters such as random drift and measurement noise have been computed. The performance of the system is evaluated for several inputs. Also, the parameter variation of output noise as a result of changing the dither characteristics is analyzed.

Keywords: Drift Error, Model Equation, Ring Laser Gyro, Scale Factor Error.

1 Introduction
Ring Laser Gyros (RLG) are effective tools for large scale geodetic surveying at a high level of accuracy. The strapdown RLG inertial systems are widely used for civilian and aircraft navigation [1-4]. They allow rotation of the sensor block with the system and analytical transformation of the output to the coordinate frame of interest, such as frequency difference between two oppositely directed laser beams. This complicated process is affected by internal and external factors, so it is difficult to measure the frequency difference and to convert it to the equivalent angle. On the other hand, test procedures vary widely depending on the facility available and the precise purpose of the test [5-6]. In this case; it is preferred to simulate the behavior of ring laser gyros. Special emphasis is directed at requirements for navigation accuracy, ambient temperature range, vibration environments, and lock in circumvention. This paper is organized as follows: In the next section, the generic model equation is described. In section 3, it is focused on the gyro modeling and characteristics determination including scale factor error, drift error and their sensitivities. In the next section, the dither influences on Gyro’s output, locking effect, scale factor and drift errors are discussed and finally the results are analyzed.

2 Gyro model Equations
All gyros have biases, scale factors, nonlinearities, noise characteristics, and so on. Many of the specification requirements and verification test methods for these parameters; their stabilities; their temperature variation, magnetic, and other sensitivities; the gyro behavior through and across vibration, and other environments; and the gyro lifetime and reliability are independent of the gyro type. The gyro model equation is defined as a series that mathematically relates its output to the components of applied gyro [7-10].

A generic model that applies to many types of sensors is shown in Figure 1. It consists of inertial (including misalignment), environmental, and random (including quantization) contributors. This approach to compartmentalizing gyro model equations is introduced to better organize the various model components. The generic model equation is

$$S_o(\Delta N/\Delta t) = (I + D + E)(1 + 10^{4\varepsilon_k})$$

where, $S_o$ is nominal scale factor, $\Delta N/\Delta t$ output pulse rate, I is the inertial input terms, E is the environmentally sensitive terms, D is the drift terms, $\varepsilon_k$ is the scale factor error terms (ppm), $\Omega$ is the equivalent gyro rate output [11] and

$$I = \omega_{irA} + \omega_{irsA} \sin \theta_N - \omega_{irsA} \sin \theta_L$$
$$E = D_\Omega \Delta T + D_{VT} \cdot V T$$
$$D = D_s + D_I(t) + D_R + D_Q$$

where

$$D_R = D_{RN} + D_{RB} + D_{RK} + D_{RR}$$

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In addition, scale factor error is defined by,

$$\epsilon_T = \epsilon_{\Delta T} \Delta T + \epsilon_{\Delta f} f(\omega)$$  \hspace{1cm} (4)$$

where, \(\omega_{\text{IRA}}, \omega_{\text{LRA}}\) and \(\omega_{\text{NRA}}\) are the components of the inertial input rate resolved into the gyro reference coordinate frame, \(\theta_L\) is the misalignment of the IA about the L reference axis (LRA), \(\theta_N\) is the misalignment of the input axis (IA) about the N reference axis (NRA). \(D_f\) is the bias, \(D_f(t)\) is the variations of \(D_f\) during warm-up period. \(D_{\Delta T}\) is the drift rate attributable to a change in temperature \(\Delta T\), where \(D_{\Delta T}\) is the drift rate temperature sensitivity coefficient. \(D_{\Delta f} \Delta f\) is the drift rate attributable to a temperature gradient \(\nabla T\), where \(D_{\Delta f} \Delta f\) is the coefficient vector of the temperature gradient drift rate sensitivity vector. 

3 Modeling

Some of the important applications of modeling occur in simulation studies, performance evaluation, and Kalman filter design. The basic difference between dynamic and stochastic modeling is as follows: in dynamic modeling, given one or more inputs (input vector), and one or more outputs (output vector), it is desired to determine the input/output relationships from both time series. Applications include those where random noise is summing at the output.

In stochastic modeling, on the other hand, there may be no direct access to an input. A model is hypothesized that, as though excited by white noise, has the same output characteristics as the unit under test. Such models are not generally unique, so certain canonical forms are chosen. For example, David Allan of the National Institute of Standards and Technology (NIST) used a power series for the PSD and the corresponding variance analysis in the time domain for the analysis of oscillator stability. This type of variance analysis is discussed in the next part. This is also the objective of the gyro drift analysis.

The general form of the model consists of a mathematical statement of the physical plant equations,
an error model consisting of a perturbation model and environmental sensitivities, a stochastic model describing random drift behavior, and a measurement model consisting of a linear combination of the output states and additive measurement noise. This is related in a generic form of a gyro model equation, consisting of the response to inertial inputs, environmental sensitivities, drift rate, and scale factor error contributors. The shortest way to gain a quantitative understanding of the RLG performance requirements is to review the RLG error sources in relation to the total system error budget [12-13].

3.1 Gyro Errors
Ring laser gyro, ideally, produce pulses representing an exact incremental change in angle. The accumulated gyro pulses, however, are corrupted by both long-term and short-term errors. Long-term errors result from instability of gyro compensation parameters and include the following terms:
- Bias error
- Scale factor error
- Input axes misalignment error

In the more demanding applications these errors may be further decomposed into temperature sensitive terms. The long -term behavior of these terms influence the error budgets established for each application and are a significant factor in designing the system.

Short-term errors cannot be calibrated and include:
- Dither spillover
- Random walk
- Quantization noise

3.1.1 Drift Errors
3.1.1.1 Random Walk Error
The basic operation of the RLG is to excite and sustain two oppositely directed traveling waves that can oscillate with different magnitudes and frequencies. The fundamental RLG equation is

$$\Delta f = \frac{4A\Omega}{L\lambda}$$  \hspace{1cm} (5)

Equation 5 shows the RLG frequency difference, commonly called beat frequency ($\Delta f$), to be proportional to the product of the geometric area enclosed by the light beams (A) and the angular rate of cavity ($\Omega_\alpha$). The beat frequency is also inversely proportional to the product of the vacuum wavelength of the laser ($\lambda$) and the optical path of cavity (L). Optical path implies sources of scale factor error. Equation 5 has been obtained with the assumption of no other sources of optical path asymmetry between the two beams other than rotation.

This frequency difference is measured optically via the two light wave interference patterns. As in any mechanical system that sustains two modes of oscillation, problems occur when the two frequencies approach each other. Energy is traded between the two modes and the frequencies tend to lock and become one. This trading of energy or coupling is largely due to the back-scattered radiation from imperfect mirrors. Loss producing mechanisms within the cavity such as outgassing of epoxies also contribute to lock-in. The major technique utilized, at this point in time to circumvent this lock-in has been called "the dither" technique. The technique consists of mechanically rocking the gyro through a stiff dither flexure suspension, which acts as a rotary spring, built into the gyro assembly. It produces a rate about the gyro input axis that causes the gyro to rapidly enter and leave the lock-in zone producing a random drift error [12, 14].

Random walk continues to be the major limitation in reducing the required time of ground gyro compassing prior to system flight. The random-walk error produces an attitude error that builds up as a function of $\sqrt{\text{time}}$. It includes four components whose coefficients are determined by Allan variance method:
- Angle random walk (Rate white noise) coefficient
- Bias instability coefficient
- Rate random walk coefficient
- Ramp coefficient

The Allan variance is a reasonable compromise. Simply put, it is a method of representing rms random drift error as a function of averaging time. It is simple to compute, much better than having a single rms drift number to apply to a system error analysis, and relatively simple to interpret and understand. It is not well suited to rigorous analysis, but a reasonable second step in the modeling process. Its most useful application is in the specification and estimation of random drift coefficients in a previously formulated model equation.

To compute these coefficients, a data set from many samples, stored during several periods, is selected. Allan variance method studies the random distributions with additive noise. In this method, the output data affected by additive input noise are predicated. It is supposed that noise distribution is predefined. Then, the noise covariance is estimated by power spectrum density.

The random walk error, $D_R$, is composed of a number of random processes of different origins. Assuming that these processes are independent, the power spectral density (PSD) of $D_R$ can be written as follows:

$$\text{PSD}(D_R) = \text{PSD}(D_{RN}) + \text{PSD}(D_{RB}) + \text{PSD}(D_{RK})$$

where, PSD(D_{RN}) is $N^2$, PSD(D_{RB}) is $B^2/2\pi f$, PSD(D_{RK}) is $K^2/(2\pi f)^2$, and $f$ is the
frequency in hertz. The coefficients N, B, K, and R can be evaluated by forming the Allan variance. Therefore,

$$\text{PSD}(D_n) = N^2 + \frac{B^2}{2\pi f} + \frac{K^2}{(2\pi f)^2} + \frac{R^2}{(2\pi f)^3}$$  \hspace{1cm} (7)$$

where, R, K, B and N are the random drift coefficients. Bias instability is considered as flicker noise in this equation. Therefore, it is necessary to calculate the temporal average of power spectrum density in Allan variance method. Figure 2 shows a scheme of the drift model [8].

From the test data after the warm-up time, compute the combined anti-lock residual and quantization noise coefficient Q by forming the following Allan variance estimates:

$$\sigma^2_{\Omega}(nT_o) = S^2[2n^2T_o^2(M - 2n)]^{-1} \sum_{m=1}^{M-2n} (N_{m+2n} - 2N_{m+n} + N_m)^2$$  \hspace{1cm} (8)$$

for n = 1, 2, 3, ..., n_{max} \leq (M–1)/2 and fitting the results to the polynomials in the least squares sense.

$$\sigma^2_{\Omega} = \frac{R^2n^2T_o^2}{2} + \frac{K^2nT_o^2}{3} + B^2\ln(2)\frac{2}{\pi}\frac{N^2}{nT_o} + \frac{3Q^2}{n^2T_o^2}$$  \hspace{1cm} (9)$$

where, S is the gyro scale factor, 1/T_o is the data sample rate, MT_o is the data record length, Nm is the total output pulses accumulated at time mT_o. The variable n should be selected smaller than the half of M. The achieved variance values will be fitted to a polynomial by using the least mean square.

### 3.1.1.2 Bias error

Gyro bias error is defined as the difference between the true low-frequency gyro bias (period greater than mission time) and the calibrated gyro bias value loaded in the computer to compensate for this error. As long as this term remains stable and the calibrated value is subtracted from the gyro output, the compensated gyro output will indicate zero rate for the zero rate input condition. Instrument factors affecting gyro bias stability are as follows:

- Stability of the mirror’s optical axis, and mirror surface.
- Outgassing of epoxy materials within the laser cavity.
- Precise path length control to correct for changes due to expansion, contraction, and bending of the gyro block material. Equation (5) shows the path length-beat frequency relationship.
- Control of the current required to sustain lasing of each beam.

Because of instability of the fixed bias at the beginning of gyro’s turn on, we can fit the resultant data to a time variant function, during the warm up time by using the least mean square method.

### 3.1.2 Scale Factor Errors

Gyro scale factor is defined as the conversion of gyro output pulse obtained by the fringe motion detector circuits into indicated angle. The basic gyro scale factor is dictated by the inertial angle change to produce one cycle of motion of the beat frequency. It directly depends on laser path length and wavelength. Also it inversely depends on the area of closed path. It is not cataloged exactly because of fabrication error, and hence, it must be measured precisely [11, 14].

Scale factor errors are as follows:

**Asymmetry:** This is a critical error source that is dramatically reduced using RLG technology over mechanical gyros. The error is defined as the unknown difference between the scale factor for positive rates and that for negative rates is given by,

$$\text{Asymmetry(PPM)} = \frac{(S_+ - S_-)}{\frac{1}{2}(S_+ + S_-)} \times 10^6$$  \hspace{1cm} (10)$$

Nonlinearity: As mentioned above, the scale factor can be obtained for any rotation rate and numerical scale factor is the mean. Maximum standard deviation indicates to nonlinear error. A simpler means of specifying nonlinearity for lesser accuracy accelerometers is to place limits on composite error as a function of input acceleration, i.e., what is commonly known as an envelope specification. The simplest envelope is to specify that the maximum error over the operating range is a specified percent of the accelerometer full-span output, especially for lesser accuracy accelerometer applications, where it is not intended to compensate for accelerometer nonlinearities that meet the specifications. Acceleration nonlinearities typically arise from structural compliance. Changes in accelerometer temperature that occur under acceleration could manifest themselves as apparent acceleration nonlinearity. The maximum of such effect could be required to be less than the nonlinearity requirement, or the heating could be regarded as a separate temperature sensitivity problem with compensation using temperature sensor measurements if possible.

Repeatability: The standard deviation of the scale factor from the mean of measurements is taken under constant conditions. Other environmental exposures may be substituted or added to the temperature environment specified.

Stability: It is the variation of scale factor between different operational cycles when the environmental conditions is stabilized. For this purpose, this sensor is set on a rotary disk with a degree of freedom and gyrates with a specific angular velocity (upper than estimated drift/bias velocity). The mean frequency is the ratio of output pulses to sampling interval. Scale factor (Hz/deg-sec or pulse/arc-sec) is obtained by using mean frequency divided by angular velocity. This experiment is repeated in several directions and finally, the mean value is considered as numerical scale factor [2]. In this work, $S_0=0.7910929$ pulse/arc-sec is obtained without considering the repeatability and stability errors of scale factor.

3.2 Gyro Sensitivities

3.2.1 Drift Sensitivities

Including:

Temperature sensitivity

Temperature gradient sensitivity

The low thermal conductivity of the RLG and its sensitivity to thermal gradients makes heating of the inertial instruments counterproductive.

Magnetic sensitivity: One of the prime mechanisms for RLG magnetic sensitivity is the property of a magnetic field to change the property of light that is not linearly polarized in such a way that a change in gyro drift occurs. All elements of the optical path are optimized to reduce this distortion of the desired linear polarized light.

3.2.2 Scale Factor Sensitivities

Temperature Sensitivity: The gyro scale factor temperature sensitivity resulting from a change in steady-state operating temperature is

$$\frac{S.F_{C_1}-S.F_{C_2}}{C_1-C_2} \times 10^6$$  \hspace{1cm} (11)

RLG technology has advanced to a point where scale factor thermal stability is good enough to make the requirement for thermal modeling unnecessary to meet thermal scale factor errors of 5 ppm.

Temperature Gradient Sensitivity: It is the change in gyro scale factor resulting from a change in steady-state temperature difference measured across the gyro.

In addition, scale factor is sensitive to the other parameters such as, voltage, temperature rate, acceleration, electromagnetic field vibration, radiance and etc [12, 16].

3.3 Measurement noise

Measurement noise is usually defined as the Allan variance component Q. The purpose of this test is to characterize the noise in the accelerometer output over a wide frequency range. Noise characterization is done by calculating the PSD, Allan variance, and/or rms deviation from the mean output in an IA vertical drift test with or without a trend removed. This noise includes two terms; the first is locking dither residual and the second is quantization noise affected by digitizing output [11, 14]. It will also be obtained by Allan variance method.

4 Results

Ring laser gyroscopes are dithered to prevent a phenomenon known as lock-in. Lock-in is the result of backscattering from the mirrors within the laser cavity causing the two counter-rotating beams to become "mode locked". This phenomenon occurs at low input rates to the gyro. By sinusoidally dithering the gyro a continuous input rate is maintained on the gyro (except for a momentary "turn-around") thus eliminating the lock-in effect. Other methods of biasing the input of an RLG have been developed. However, only the dither method has been
successfully transitioned to production [17]. Therefore, we use a sinusoidal dither with random amplitude which eliminates the locking interval (Figure 3). On the other hand, it increases the angle random walk noise considerably, confined to quantum values by using of the anti-vibrating structures of gyro’s frame in the ideal mode. Figure 4 shows the angle random walk noise as a function of cavity length for different laser output power.

As shown in Figure 4, the angle random walk decreases with increasing the cavity length and it increases with increasing the laser output power. In this work, the angle random walk noise is obtained as $0.0023 \times \sqrt{n}$ which is comparable to available test data in [12, 18].

The dither contributes to nonlinear error of scale factor, related to input variations. Figure 5 shows a plot of the dithered scale factor correction, expressed in parts per million, for dither rates of 100, 200, 300 deg/sec. Both the positive and negative contributions are plotted on the same scale. The negative term has been multiplied by -1, so that it can be scaled relative to the positive term. The value of $\Omega$ was chosen so that the maximum values of the two contributions are approximately equal [9]. The result revealed when the rotating rate is equal to the peak dither rate, the error is maximal. The reason is that the interval locking of sensor is increased to maximum value. Here, the maximum error is 250 ppm.

In addition, the scale factor error is affected by cavity gain and loss (Figure 6). Instability of bias is increased since the mean value of dither rotation is not zero during the sampling period. This parameter is 0.003deg/h.

Furthermore, the offset created by the current difference of laser, discharges in two directions, and increases instability of bias. Figure 7 shows the bias discharge current as a function of total discharge. The response related to the sinusoidal input with frequency of 0.1 Hz and amplitude of 10 deg/sec is depicted in Figure 8(a). Also the response related to the fixed input with the peak dither rate of 174 deg/sec, is shown in Figure 8(b).

**Fig. 3 Dither scheme [9].**

**Fig. 4 Angle random walk noise versus cavity length.**
Fig. 5 These theoretical plots show the scale factor deviation from an arbitrary nominal, for three peak dither rates. The parameters have been adjusted, so as to overlap the positive and negative SF corrections. The negative SF correction shows a smaller width in rotation space, than the positive correction, agreeing with experimental results. The calculated expression, for the peak SF correction when the rotation rate is equal to the dither rate is also shown. Note that it agrees well with the values of the peaks shown on the plot [9].

Fig. 6 Scale factor error as a function of gain and loss.
Bias-Discharge Current Diagram

Fig. 7 Bias discharge current versus total discharge current.

Fig. 8 (a) The response related to the sinusoid input with frequency of 0.1 Hz and amplitude of 10deg/sec (b) The response related to the fixed input with the peak amplitude of rotation rate.

5 Conclusion
In this work, the RLG system with respect to its performance has been addressed. It is modeled by using two methods; in the dynamic model, some parameters such as scale factor and environmental sensitivity have been determined, whereas in the stochastic model, the other parameters such as random drift and measurement noise have been computed. The results revealed that the most effective errors found to be the bias instability and the angle random walk increased by using the dither technique. However these errors will be decreased through the judicious choice of frequency, peak rotation rate and eliminating the temporal mean of rotation rate.

References


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Shahram Mohammad-Nejad received his B.Sc in Electrical Engineering from University of Houston, Houston, USA, in 1981 and M.S. and Ph.D. degrees in Semiconductor Material Growth and Lasers from Shizuoka University, Shizuoka, Japan, in 1990 and 1993, respectively. Professor Mohammad Nejad invented the PdSrS laser for the first time in 1992. He has published over 80 scientific papers and books. Currently, he is the Head of Electrical Engineering Department, Iran University of Science and Technology, Tehran, Iran. Also, he is a scientific committee member of Iranian Conference on Electrical Engineering (ICEE) and also Iranian Conference on Optics & Photonics, member of Institute of Engineering and Technology (IET) and an IET-CEng. His research interests include semiconductor material growth, quantum electronics, semiconductor devices, optoelectronics, electronic devices and lasers.

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