

# Determination of Electric Field at Inception Based upon Current-Voltage Characteristics of AC Corona in Rod-Plane Gaps

H. Javadi\*, M. Farzaneh\*\* and A. Peyda\*\*

**Abstract:** This paper deals with the measurement of AC corona inception voltage,  $V_{incp}$ , at the tip of a rod electrode using a hemispherically-capped rod-plane electrode configuration for various rod radii with a short air gap. Effects of atmospheric pressure and temperature variation on  $V_{incp}$  are investigated experimentally. An empirical equation for the field form factors of the hemispherically capped rod-plane electrodes is proposed with its range of applicability. The obtained results are analyzed to derive a more accurate analytical equation for the calculation of the electric field at corona inception voltage,  $E_{incp}$ , and the average of electric field distribution,  $E_{mean}$ .

**Keywords:** Electric field, Corona Inception Voltage, Rod-plane Electrodes, Form Factor, Atmospheric Correction Factor, 3-D Electric Field Simulation.

## 1 Introduction

MANY investigations have been carried out to observe and understand the complex processes underlying breakdown across air gaps. These studies include: (i) the physical process of corona and breakdown across air gaps; (ii) the electric field calculation for various electrode configurations; and (iii) the effect of atmospheric parameters.

Studies in the first category aim at understanding the phenomena and providing explanations of the microscopic interactions involved (electron attachments, photo ionization, velocity of charged particles,...), as well as their manifestations in partial and various arc forms and avalanches, and the relation of discharge current to the different contributing voltages [1-6]. In most of these studies, dc or impulse voltage was used for medium to long air gaps. Although this research provides a good insight into the physics of the breakdown process, it barely has pointed to any clear computation technique or proposed adequate equations

for the electric field at the corona inception voltage,  $E_{incp}$ .

For the second category of investigations, about electric field calculation for various electrode configurations, most of the literature deals with the numerical computation of electric fields using computer modeling such as the Finite Element Method "FEM" [7] and the Charge Simulation Method "CSM" [8-9] which, though unable to derive analytical formulae, provide close computer simulation. Also, it includes an expression derived for short air gaps in a rod-plane configuration with SF<sub>6</sub> insulating medium [10]. However, all these methods provide results which are not validated with laboratory tests for different voltage wave shapes.

There are also a few useful empirical formulae for other highly non-uniform electrode configurations [11-12], but there is no equation for the calculation of  $E_{incp}$  at the tip of the rod in a rod-plane electrode system. The only known equation for that purpose was derived by Coelho [13] using a hyperbolic approximation, and has been used in a variety of applications [14-17]. As it is known from the extensive research on the breakdown process, there are different stages of the development of streamers and partial discharges, each depending on a different set of factors and degrees of contribution to the evolving non-uniformity from inception of streamer corona to complete breakdown. It is almost impossible to explain the breakdown process generally for a vast range of rod-plane gaps, because the streamer build-up process varies among a set of gap and rod radius sizes.

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Therefore, one should not assume that this equation is adequate for all rod-gap size ranges. However, some empirical expressions were derived to calculate field form factors in order to estimate the values of the ratio  $E_{max}/E_{mean}$  [18-19] for practical purposes in the case of a rod-plane electrode system. The main issue here is that the average value of the electric field distribution,  $E_{mean}$ , cannot be adopted universally as it depends on electrode geometry and on the wave shape of the applied voltage.

In the third category of investigations, dealing with the effects of atmospheric parameter variations on the breakdown voltage of rod-plane gaps [21-22], all the studies have shown that the IEC correction method [23-25] has many setbacks and cannot be generally applied to all electrode configurations and voltage wave shapes. Just a few investigations have been carried out on the effect of humidity for a 20 cm air gap, showing that there is a large difference between the threshold and stability fields for smaller air gaps.

The setback arises from the fact that the effect of air density seems to be dependent on the combination of air pressure and temperature, suggesting that the only feasible way to perform tests is to keep one parameter constant while varying the other. Furthermore, previous studies have shown the necessity of further research on the effect of temperature and pressure variation to determine their interaction on AC corona inception voltage,  $V_{inc}$ .

## 2 Experimental Facilities

Fig. 1 shows the schematic diagram of the experimental setup. The high voltage ac transformer of the setup provides the necessary power and voltage to the circuit [26, 27]. The voltage is regulated by a control unit connected to a stabilizing capacitor  $C_R$ . The circuit consists of a current limiting resistor,  $R_d$ , and a resistive voltage divider comprises two series resistors,  $R_{m1}$  and  $R_{m2}$ , where these resistors act respectively as the HV arm of the voltage divider and the LV resistor of the measuring arm.

A custom-built interconnecting bushing connects the high voltage circuit to the test chamber through the Faraday cage. The plane electrode is attached to the common ground wire and to the leakage current sensing resistance  $R_c$ . The voltage and current signals are fed to a data acquisition system through a voltage conditioner/limiter. The voltage limiter's cut threshold is  $\pm 10V$  (peak). The sampling rate is set to 800 samples with a 4000Hz clock rate. The data is measured and stored using the LABVIEW<sup>®</sup> software.

To keep the corona inception voltage measurement as high as accurate as possible, humidity, air temperature and air pressure were kept nearly constant for each set of tests. In order to ensure precision of the test results, a standard system (sphere-sphere electrode configuration, Fig. 2) was used to measure the breakdown voltage at the start of each test set. This value was subsequently used to convert the obtained

corona inception voltages for that particular test to standard values.

Fig. 3 shows the set of rod electrodes used in the test object. With the available facilities, rod electrodes of radii have been produced by an acceptable level of accuracy. The radii of the rod electrodes were ranged between 0.25 mm and 3.04 mm, with gap spacing ranging from 1 cm to 15 cm in a 1-cm step increment. In other words, this standard geometry delivers a good reference point for converting the obtained values to the standard atmospheric conditions ( $T=20^\circ C$ ,  $P=101.3kPa$  and  $h=11g/m^3$ ).

The peak value of applied voltage at the standard conditions is 31.8kV for the specific electrode system used, thus yielding an electric field of  $31.8kVcm^{-1}$ , which is the minimum required to start corona discharge under the test conditions.

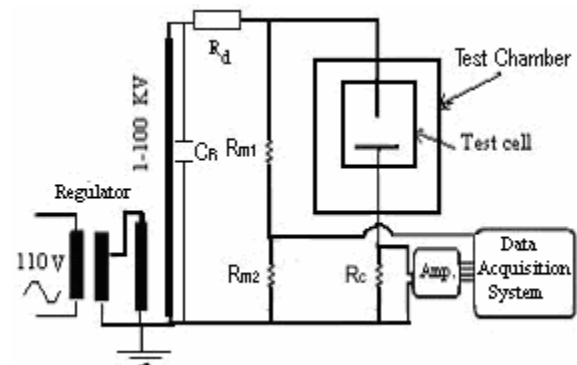


Fig. 1 Circuit diagram of testing setup.



Fig. 2 Standard sphere-sphere electrode.

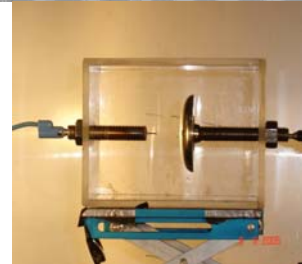
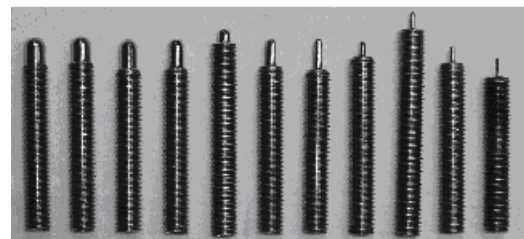


Fig. 3 Hemispherically-capped steel rod electrodes and test object.

Considering the statistical nature of breakdown, in each case,  $V_{incp}$  was measured five times and the mean value was recorded.  $V_{incp}$  is easily detectable because of the appearance of partial discharge, and the deformation of the current wave shape including a phase shift from 90 to 0 degree (pure capacitive current towards pure resistive one). In order to measure the corona inception voltage, three traditional methods were evaluated.

1. Audio detection
2. Visual detection
3. Current wave shape

So, the analysis of current wave shapes was found to be the best way to measure the corona inception voltage. Current and voltage were simultaneously recorded with a voltage rise rate of 5kV/sec. Since partial discharges appear differently on the current waveforms for different radii and electrode gap distances, these recordings were thoroughly studied to figure out a suitable protocol for corona inception detection by discharge current waveform.

The maximum variations of the atmospheric parameters have been recorded for each set of tests, and were found to be:

1. Ambient temperature: 4°C
2. Atmospheric pressure: 0.005b
3. Humidity: 8 % = 1.56gr/m<sup>3</sup> (this humidity change occurred between two extreme values of 7.7 and 9.24gr/m<sup>3</sup>)

It should be noted that these maximum variations did not occur during the same test sets and did not cause any error in the results.

### 3 Experimental Results

The results of the laboratory tests are subdivided into three categories. In the first category, we find the measured values of corona inception voltage for the purpose of deriving an equation to calculate  $E_{incp}$  and  $E_{mean}$  at the tip of the rod. The second category comprises the data obtained from atmospheric pressure tests. In the third category the last measured values of corona inception voltage under temperature variation are found.

#### 3.1 Measured $V_{incp}$ using the Rod-Plane Electrode Configuration

Considering the unit procedure for the whole geometry range ( $0.25\text{ mm} < r < 3.04\text{ mm}$  and  $1\text{ cm} < d < 15\text{ cm}$ ), leakage current and voltage across the gap were recorded using the LabView software. Since peculiar current wave shape deformities are associated with different gap geometries, all the waveforms were meticulously analyzed and studied to detect the point of corona inception.

Fig. 4 shows the detection point of corona inception voltage in a rod-plane electrode configuration with  $r=2.48\text{ mm}$  and  $d=4\text{ cm}$ . Generally, this point is characterized by a slight degree of deformity, an

envelope of partial discharges appearing as noise on the gap leakage current waveform and departure of current phase from 90 to 0 degree.

During the laboratory tests, the corona inception regions were observed. From the experimental results, the following categories of corona inception have been distinguished:

- For large-diameter rod tips, no corona was observed during the whole breakdown phenomena. Instead, the initial ionization process manifested itself in the form of disruptive streamer formations. One of the main reasons for that is that the corona critical length is longer than the gap spacing.

- For  $0.5\text{ mm} < r < 2.58\text{ mm}$  values and for the whole range of gap spacing, a streamer corona developed at an early stage, with a corona film appearing at a later stage.

- For  $0.25\text{ mm} < r < 0.5\text{ mm}$  and  $1\text{ cm} < d < 4\text{ cm}$  values, there was a very fast transition period from corona to leader development.

- For  $0.25\text{ mm} < r < 1\text{ mm}$  and  $5\text{ cm} < d < 10\text{ cm}$  values, there was a clear and visible presence of corona film on the tip of the rod during a noticeable period.

Fig. 5 also shows a plot of the measured  $V_{incp}$  for the whole range of radii  $0.25\text{ mm} < r < 3.04\text{ mm}$ . As can be noticed, the plot tends to follow a saturation curve pattern.

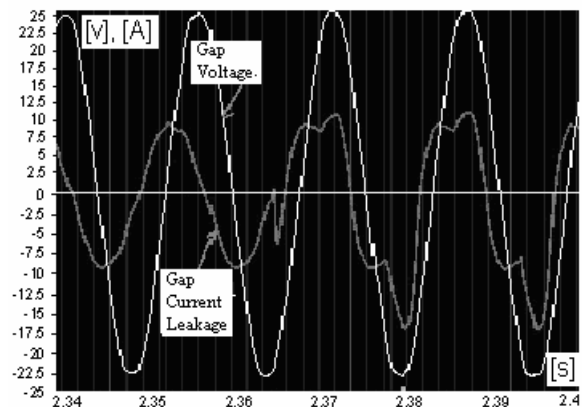


Fig. 4 Corona inception voltage and current wave shape for a rod-plane gap with  $r=2.48\text{ mm}$  and  $d=4\text{ cm}$ .

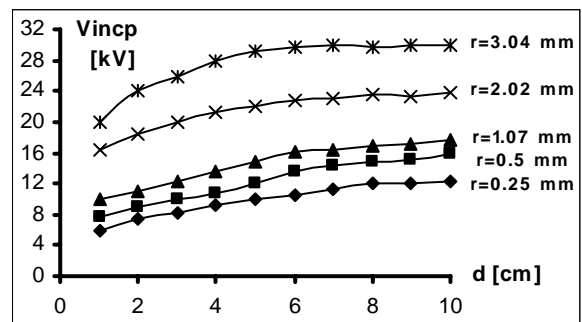


Fig. 5 Corona inception voltage of the rod-plane gaps vs. gap spacing.

### 3.2 Atmospheric Temperatures Tests

Using a climatic chamber,  $V_{incp}$  was measured for a temperature range of  $-20^{\circ}\text{C} < T_{amb} < 20^{\circ}\text{C}$ . The results are plotted in Fig. 6, where the measured values are depicted as circles, and where the dotted line represents the result of a linear regression on these values representing their power trend line. The geometry for these atmospheric tests was set to  $r=1.54\text{mm}$  with a gap of  $d=7\text{cm}$ , which approximately represents the mean value of the experimental setup geometry. The plot reveals that a temperature decrease from room temperature to  $-20^{\circ}\text{C}$  has a relatively slight effect on  $V_{incp}$ . Over that 40-degree range,  $V_{incp}$  increases by about 3.4kV. The test data for the temperature tests are displayed in Table 1. During the tests, the monitored parameters were pressure, temperature, dew point (DP) and relative humidity (RH).

### 3.3 Atmospheric Pressure Tests

The measured  $V_{incp}$  for an ambient pressure,  $P_{amb}$ , ranging from 22.8kPa to 101.8kPa, which roughly corresponds to an altitude ranging from 10,500m to 30m, is plotted in Fig. 7, where the circles represent the measured values and where the dotted line is a power curve, fitted to these. This plot shows the fact that air pressure variation affects corona inception voltage more significantly than variation in ambient temperature. For the given pressure range,  $V_{incp}$  decreases as low as 9kV at 22.8kPa, which is a considerable drop as compared to that of 21.2kV at 101.8kPa (the altitude of the laboratory in Chicoutimi is roughly 30m above sea level). The chosen geometry for the atmospheric tests is the same as that used for the ambient temperature tests.

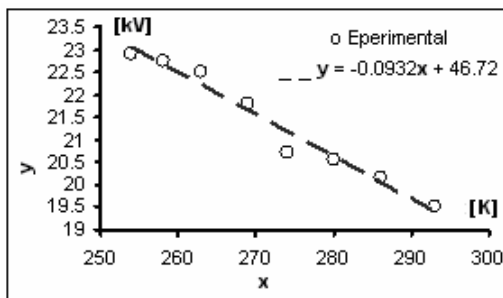


Fig. 6 Effect of decrease in ambient temperature on corona inception voltage for a rod-plane configuration with  $r=1.54\text{mm}$  and  $d=7\text{cm}$ .

Table 1 Monitored ambient parameters during the temperature tests.

Air Pressure=101.07 kPa							
Air Temperature °C							
°C	20	13.7	6.8	0.4	-10	-15	-19.4
$V_{incp}$ (kV)	19.5	20.1	20.6	20.7	22.5	22.8	22.9
$h(\text{gr}/\text{m}^3)$	4.1	3	2.2	1.6	0.6	0.3	0.3
RH (%)	24	25.4	29.5	32.9	37	43.6	53
DP (°C)	-1.3	-5.9	-10.1	-14.7	-23.5	-24.7	-26.3

Table 2 provides the monitored ambient parameters during the tests.

## 4 Computer Simulation

Since electric field is not measurable experimentally, computer simulation is needed to determine  $E_{max}$  and the form factor of each rod-plane electrode.

This simulation was carried out using the Finite Element Method (FEM) based on the Coulomb-3D software for solving static field problems in a three dimensional medium.

### 4.1 Modeling of Rod-Plane Gaps

For the purpose of computer modeling, exact replicas of the physical configurations were used. The all boundaries of study area were considered the Neumann's type except at the rod electrode (as high voltage electrode using 1V in per unit value) and at the plane electrode (as ground electrode using 0V) considered as Dirichlet's type. Fig. 8 depicts a 3-D model for a rod-plane gap with  $r=1.54\text{mm}$  and  $d=4\text{cm}$ , used for atmospheric temperature and pressure tests. Fig. 9 illustrates the distribution of electric strength on a profile plane perpendicular to the  $y$ -axis for the same gap arrangement. In that figure, the magnitude of the electric strength is represented by color shades. As expected, the maximum field strength,  $E_{max}$ , in the air gap is located at the tip of the rod, right on the central longitudinal axis of the rod.

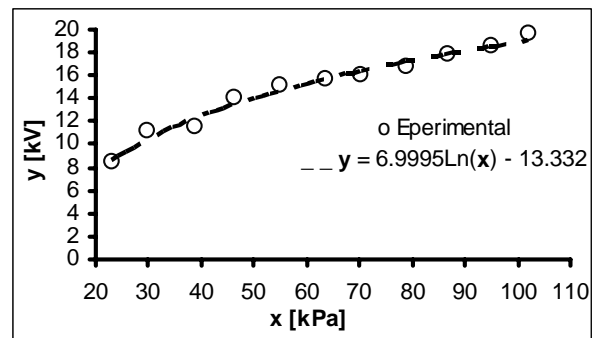
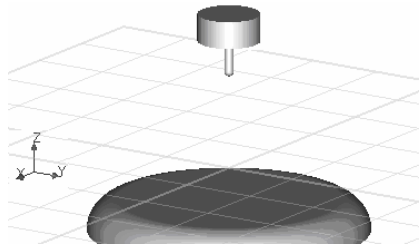


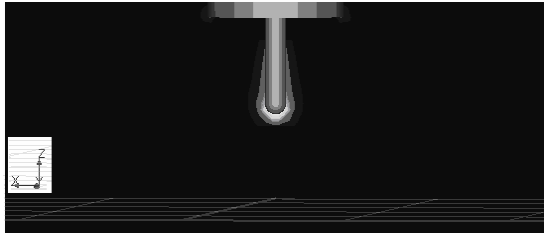
Fig. 7 Effect of ambient pressure reduction on the corona inception voltage for a rod-plane arrangement with  $r=1.54\text{mm}$  and  $d=7\text{cm}$ .

Table 2 Monitored ambient parameters during the pressure tests.

Air Temperature=21°C							
Air Pressure kPa							
kPa	101.8	94.8	86.8	78.8	54.8	38.8	22.8
$V_{incp}$ (kV)	19.5	18.5	17.8	16.7	15	11.5	8.3
$h(\text{gr}/\text{m}^3)$	3.42	3.24	3.2	3.06	2.53	2.1	1.9
RH (%)	18.3	17.2	17.2	16.6	14.7	12.6	11.6
DP (°C)	-3.7	-4.4	-4.6	-5.2	-6.8	-8.9	22.8



**Fig. 8** 3-D model of a rod-plane gap of  $r=1.54\text{mm}$  and  $d=4\text{cm}$ .



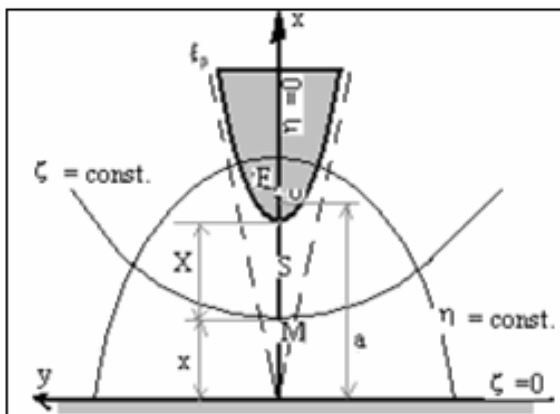
**Fig. 9** Distribution of electric field in the gap area for an arrangement with  $r=1.54\text{mm}$  and  $d=4\text{cm}$ .

#### 4.2. Analytic and Simulation Results

*Coelho et al.* proposed an equation based on a simple hyperboloid geometric approximation to calculate the electric field in a point-plane gap. A tip can be generated by the hyperbola of Eq. (1) rotating around the  $x$  axis shown in Fig. 10.

$$\begin{aligned} x &= a \sin \xi \cosh \eta \\ y &= a \cos \xi \sinh \eta \end{aligned} \quad (1)$$

where  $\zeta$  is the parameter defining the hyperbola. If  $\zeta=0$ , it corresponds to the  $Oy$  axis generating plane. If  $\zeta=\pi/2$ , ( $y=0$ ), it corresponds to the  $Ox$  axis generating an infinitely sharp hyperbola.  $\eta$  is the parameter defining a particular point on the hyperbola defined by  $\zeta$ . The value of the electric field at point  $(\zeta, \eta)$  is given by Eq. (2):



**Fig. 10** Representation of a point-plane configuration with the various quantities involved.

$$E(\zeta, \eta) = \frac{C}{a \cos \zeta (\cosh^2 \eta - \sin^2 \zeta)^{\frac{1}{2}}} \quad (2)$$

where  $C$  is a constant depending on the applied voltage. If we set  $\eta=0$ , the above equation then represents the field along the  $x$  axis. In this particular case, making use of some approximations, at the tip ( $X=0$ ), Eq. (3) is obtained:

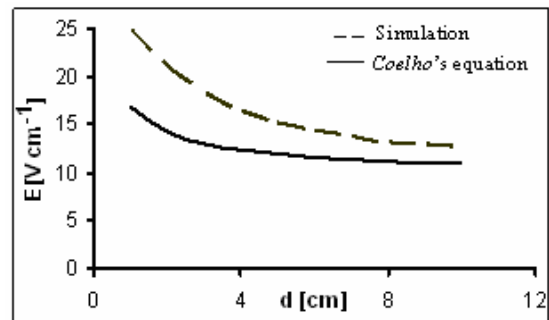
$$E(0) = \frac{C}{r} = \frac{V}{r \ln \left( 2 \left( \frac{a}{r} \right)^{\frac{1}{2}} \right)} \quad (3)$$

Assuming that  $d \approx a$  in Eq. (3), the electric field strength at the tip of the rod can be analytically obtained using Eq. (4).

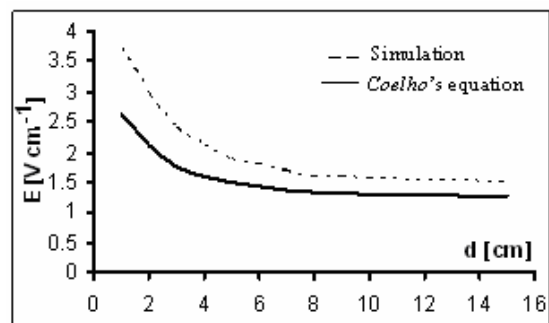
$$E_{incp} = \frac{2V}{r \ln \left( \frac{4d}{r} \right)} \quad (4)$$

where  $V$  is the applied voltage,  $r$  is the rod radius, and  $d$  is the gap length.

In Fig. 11, the simulation results and the values obtained by *Coelho's* equation are compared. Both plots represent  $E_{max}$  at the tip of the rod for a gap distance of  $1\text{cm} < d < 10\text{cm}$  for  $1\text{cm}$  increments, and a distance of  $15\text{cm}$  in the case of an applied voltage of  $1\text{V}$  on the rod electrode.



(a)



(b)

**Fig. 11** Electric field at the tip of the rods vs. gaps for an applied voltage of  $1\text{V}$  on the rod. (a)  $r=0.25\text{mm}$  (b)  $r=3.04\text{mm}$

It is obvious from these plots that:

- The smaller the radius of the rod and the shorter the gap, the higher is the error associated with *Coelho's* equation.
- The maximum error of *Coelho's* equation occurs at  $r=0.25\text{mm}$  and  $d=1\text{cm}$  with the value of  $9\text{Vcm}^{-1}$ . The minimum error occurs at  $r=3.04\text{mm}$  and  $d=15\text{cm}$  with the value of  $0.25\text{Vcm}^{-1}$ .
- *Coelho's* equation provides fairly good results for  $r>1.54\text{mm}$  and  $d>10\text{cm}$ .

## 5 Equation and Data Analysis

The test and simulation results were used to derive equations to calculate  $f$ ,  $E_{mean}$  and  $E_{incp}$  for the prescribed ranges of the rod-plane electrode arrangements. Moreover, correction factors for standard atmospheric conditions for air pressure and temperatures were also derived from the experimental results.

### 5.1 Field form Factor Expression

Using the values of  $E_{max}$  and  $E_{mean}$  following the computer simulations, the field form factors were calculated by means of Eq. 5,

$$f = \frac{1}{\gamma} = \frac{E_{max}}{E_{mean}} \quad (5)$$

where  $E_{mean}$  is simply the ratio of the corona inception voltage of a rod-plane gap to the gap spacing. The results are demonstrated in Fig. 12.

Then, using Matlab's curve fitting tool package, the best fit was applied to the obtained data to achieve an analytic relation. Equation 6 is derived from the data obtained by simulation in order to calculate form factors for the prescribed gap to radius range.

$$f = n.e^{\left(\frac{j.d}{r}\right)} + s.e^{\left(\frac{l.d}{r}\right)} \quad \text{for } 3 \leq \frac{d}{r} \leq 400 \quad (6)$$

where  $n=60.79$ ,  $j=0.002$ ,  $s=-58.1$ , and  $l=-0.006$ .

### 5.2 Expression for Calculating Mean Electric Strength

Normally for an electrode configuration, typically a parallel plane configuration, the point of breakdown and corona discharge inception voltages are the same, i.e. around  $31\text{ kVcm}^{-1}$  in the air. Nevertheless, in the case of configuration producing highly non-uniform electric fields around their gap, this value varies considerably. Some factors affect that value are atmospheric conditions, electrode shape and arrangement and applied voltage waveform.

It would be very useful for researchers or engineers to have at their disposal an expression to calculate the mean electric strength in a rod-plane air gap at the point of the corona/PD inception (AC voltage), thus avoiding

the use of expensive test equipment and time-consuming computer simulation.

In order to derive such an expression, after plotting the  $V_{incp}$  values obtained by laboratory tests against the gap spacing, regression analysis was applied to obtain the best fit trend line for all the plots, yielding Eq. (7):

$$E_{mean} = Ad^B \quad (7)$$

where coefficients  $A$  and  $B$  are fitting constants for each trend line, and  $d$  is the gap spacing.

The second step is to plot the constants,  $A$  and  $B$ , against the radii, with a resulting linear regression for  $A$  and a logarithmic one for  $B$  against  $r$ , as shown in Eq. (8).

$$A = k_1.r + k_2 \quad \text{and} \quad B = k_3.\ln(r) + k_4 \quad (8)$$

where  $k_1$  to  $k_4$  are known coefficients. By substituting the  $A$  and  $B$  values from Eq. (8) into Eq. (7), we obtain:

$$E_{mean} = \frac{4.3103 \times r + 5.8558}{d^{(0.021\ln(r)+0.781)}} \quad (9)$$

In this equation,  $E_{mean}$  is in  $\text{kVcm}^{-1}$ ,  $r$  in mm and  $d$  in cm.

Fig. 13 shows the plots of Eq. (9) against gap spacing for two different rod radii. It also compares the values obtained by that equation with experimental ones.

Table 3 also shows the maximum and minimum deviations of the derived equation from the experimental values in percentage as defined in Eq. (10).

$$Dev = \left| \left( \frac{E_{mean(eq_n)} - E_{mean(exp.)}}{E_{mean(exp.)}} \right) \times 100 \right| \quad (10)$$

### 5.3 Calculation of $E_{incp}$

The corona onset field strength at the tip of the rod can be calculated from a combination of experimental and computer simulation results. This is achieved by Eq. (11).

$$E_{incp} = E_{mean} \times f = \frac{V_{incp} \times f}{d} \quad (11)$$

where  $E_{incp}$  is in  $\text{kVcm}^{-1}$ ,  $E_{mean}$  is in  $\text{kVcm}^{-1}$  and  $f$  will be dimensionless.

For that purpose, the same procedure to derive  $E_{mean}$  is used to calculate the maximum electric strength at the tip of the rod at the corona inception point, yielding Eq. (12):

$$E_{incp} = \frac{89.576}{d^{(0.035\ln(r)+0.091)} \times r^{0.365}} \quad (12)$$

For  $0.25\text{mm} \leq r \leq 3.04\text{mm}$  and  $1\text{cm} \leq d \leq 15\text{cm}$

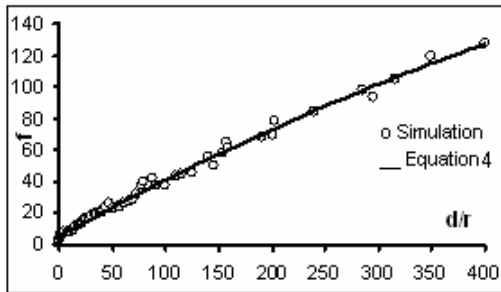


where  $E_{incp}$  is in  $kVcm^{-1}$ ,  $r$  is the rod radius in mm and  $d$  is the gap length in cm. Table 4 shows maximum and minimum deviations of the values obtained by Eq. (12) from the experimental results.

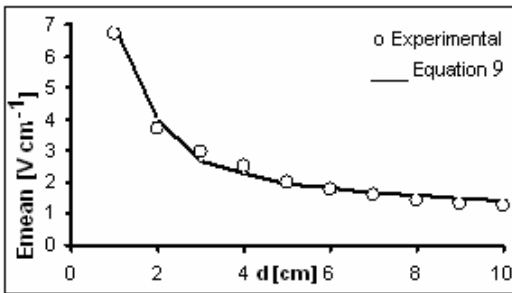
Also in order to verify Eq. (12), Fig. 14 and Fig. 15 show respectively the results obtained from that equation and those from the hyperboloid method (Eq. 4).

**Table 3** Deviation of Eq. (9) from the Experimental Values in Percentage.

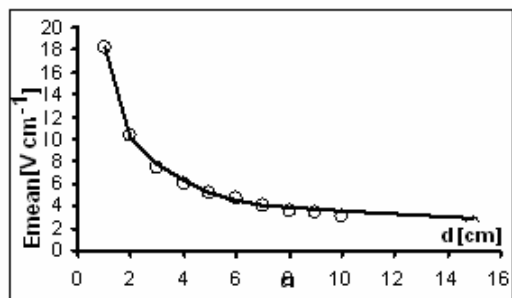
Maximum & Minimum deviations (%)			
	Rod Radius (mm)	Gap (cm)	Deviation
<b>Max</b>	1.07	15	14.09%
<b>Min</b>	3.04	8	0.02%



**Fig. 12** Form factor value as the function of  $d/r$ .



(a)

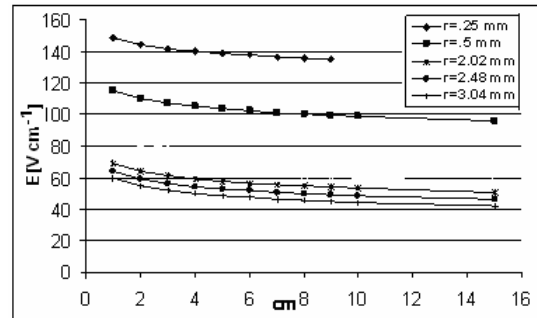


(b)

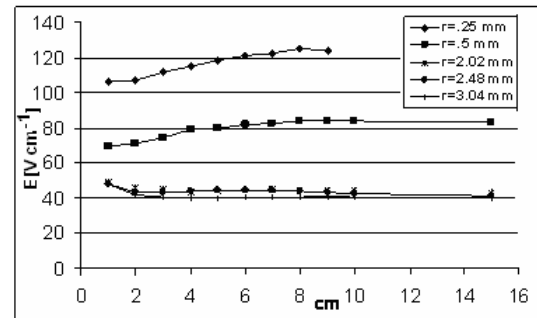
**Fig. 13**  $E_{mean}$  as a function of gap spacing for two different radii at the corona inception point. (a)  $r=0.25mm$  (b)  $r=3.04mm$

**Table 4** Deviation of Value Determined by Eq. (12) as Compared to Experimental One.

Maximum & Minimum Deviations (%)			
	Rod Radius (mm)	Gap (cm)	Deviation
<b>Max</b>	1.07	15	19%
<b>Min</b>	0.5	8	0.11%



**Fig. 14** Electric field at the tip of the rod, at the corona inception obtained by Equation 12 - vs. gap distance.



**Fig. 15** Electric field at the tip of the rod, at the corona inception obtained from Equation 4 - vs. gap distance.

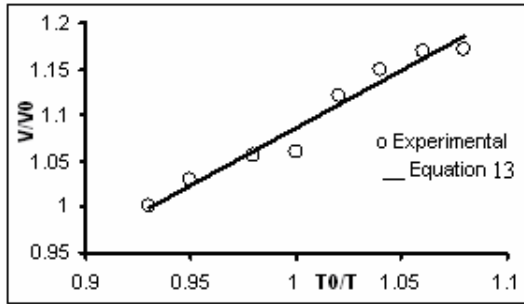
## 6 Correction Factors for Atmospheric Conditions

In order to achieve a precise temperature correction factor, the same regression method as the one used to derive  $E_{mean}$  (Eq. 7) was applied to the experimental data, thus yielding Eq. (13):

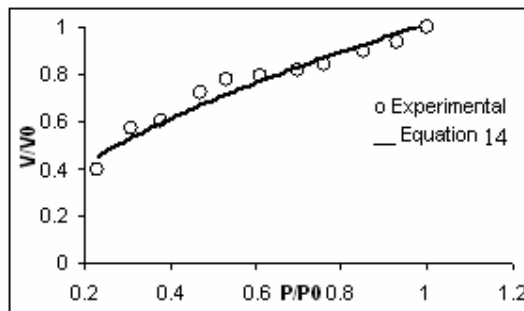
$$\frac{V_{incp}}{V_{incp(0)}} = 1.09 \left( \frac{T_0}{T} \right)^{1.19}; 253^\circ K < T < 293^\circ K \quad (13)$$

where  $V_{incp}$  stands for corona inception voltage at temperature  $T$  in kV (peak),  $V_{incp(0)}$  is the inception voltage at the atmospheric standard condition, and  $T_0$  is the standard atmospheric temperature, equal to  $20^\circ C$  or  $293^\circ K$ . That equation is plotted in Fig. 16 against the experimental results, with a maximum error of 8%.

In order to determine the correction factor due to pressure, the power regression method was applied to the  $(V_{incp}/V_{incp(0)})$  and  $(P_{amb}/P_{amb(0)})$  ratios of the experimental data sets, yielding Eq. (14), which is the analytical form for the calculation of  $V_{incp}$  for the standard conditions:



**Fig. 16** Ratio of  $V/V_0$  against ratio of  $T_0/T$  for  $r=1.54\text{mm}$  and  $d=7\text{cm}$ .  $T_0=273\text{ °K}$  and  $V_0=19.5\text{ kV}$  (peak).



**Fig. 17** Ratio of  $V/V_0$  against ratio of  $P/P_0$  for  $r=1.54\text{mm}$  and  $d=7\text{cm}$ .  $P_0=101.3\text{ kPa}$  and  $V_0=19.5\text{ kV}$  (peak).

$$\frac{V_{incp}}{V_{incp(0)}} = \left( \frac{P_{amb}}{P_{amb(0)}} \right)^{0.5239}; 22.8\text{ kPa} < P_{amb} < 101.8\text{ kPa} \quad (14)$$

where  $V_{incp}$  stands for the corona inception voltage at the ambient pressure,  $P_{amb}$ , in kV (peak), and where  $V_{incp(0)}$  is the inception voltage at the atmospheric standard conditions, and  $P_{amb(0)}$  is the standard atmospheric pressure, equal to 101.3 kPa. Fig. 17 shows the results obtained from Eq. (14) against the experimental ones.

## 7 Conclusion

Using a rod-plane electrode configuration with a radius range of  $0.25\text{mm} < r < 3.04\text{mm}$  and a gap spacing range of  $1\text{cm} < d < 15\text{cm}$ , a series of tests were carried out to measure the corona inception voltage. Computer simulation was carried out to obtain the field form factors for the electrode configuration used in the experimental setup. Also, the effects of atmospheric pressure and temperature were investigated on an electrode configuration of  $r=1.54\text{mm}$  and  $d=7\text{cm}$  with air temperature variation range of  $-20\text{ °C} < T_{amb} < 20\text{ °C}$  and ambient pressure variation range of  $22.8\text{ kPa} < P_{amb} < 101.8\text{ kPa}$ .

Empirical equations were derived from the experimental and simulation results to calculate  $E_{mean}$ ,  $f$  and  $E_{incp}$ . Moreover, correction factors for the atmospheric parameters were derived from the experimental results. The results and observations are summarized as follows:

1. Computer simulation results reveal that Eq. (13) generates very low values of  $E_{max}$  particularly for a rod-plane electrode configuration with  $r < 1.5\text{mm}$  and  $d=4\text{cm}$ . The maximum error using *Coelho's* equation occurs at  $r=0.25\text{mm}$  and  $d=1\text{cm}$  with a value of  $9\text{Vcm}^{-1}$ , while the minimum error occurs at  $r=3.04\text{mm}$  and  $d=15\text{cm}$  with a value of  $0.25\text{Vcm}^{-1}$ . It appears from the simulation results that *Coelho's* equation returns fairly good results for the rod-plane configuration with  $r > 1.54\text{mm}$  and  $d > 10\text{cm}$ .

2. An empirical expression, Eq. (6), for calculating the field form factors for the experimental setup was derived. This equation will only be valid for  $3 < d/r < 400$  within the prescribed electrode arrangement. The results show that for this particular range, that value may reach up to 125 for  $(d/r)=400$ .

3. Based on the experimental results, Eq. (9) can be used to calculate  $E_{mean}$  for the setup electrode arrangements when the applied voltage is  $V_{incp}$ . This expression has a maximum error of 14%.

Finally, Eq. (12) was derived from both experimental and simulation results for calculating  $E_{incp}$  at the tip of the rod in a rod-plane gap. This equation provides accurate and realistic results which can replace the use of *Coelho's* equation for the prescribed rod-plane geometry ranges.

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